

Nature convective et diffusive de la méthode de Boltzmann sur réseau

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- 1 Éléments fondateurs de la LBM
 - Liens avec l'équation de Boltzmann continue
 - Les vitesses discrètes
 - L'algorithme LBM
- 2 Analyse convective du modèle
 - Modes du modèle NS compressible isotherme
 - Modes LBM
 - Selective and adaptive filtering
 - Adaptive relaxation times
 - Quelques applications
- 3 Interfaces et Mélanges complexes
 - Interfacial dynamics
 - Multicomponent mixing

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Equation de Boltzmann (1872)

$$\frac{\partial f}{\partial t} + c_i \frac{\partial f}{\partial x_i} + \frac{F_i}{m} \frac{\partial f}{\partial c_i} = \left(\frac{\partial f}{\partial t} \right)_{coll} \quad (1)$$

$f(\mathbf{x}, \mathbf{c}, t)$ Densité de probabilité qu'une particule de vitesse $\mathbf{c}$ soit en (x_1, x_2, x_3) au temps t .



Les quantités macroscopiques $\rho, \rho U, \rho e$ sont obtenus par les moments :

$$\int (.) d\mathbf{c}, \quad \int c (.) d\mathbf{c}, \quad \int c^2 (.) d\mathbf{c}$$

En 1954, P.,L. Bhatnagar, E.P. Gross et M. Krook montrent que l'EDP :

$$\frac{\partial f}{\partial t} + c_i \frac{\partial f}{\partial x_i} = -\frac{1}{\tau} [f - f^{eq}] \quad (2)$$

Avec la fonction d'équilibre :

$$f^{eq}(\mathbf{x}, \mathbf{c}, t) = \rho \left(\frac{m}{2\pi k_B T} \right)^{3/2} \exp \left[-\frac{m}{2k_B T} (\mathbf{c} - \mathbf{u})^2 \right] \quad (3)$$

Est équivalent à Navier-Stokes pour les faibles nombres de Knudsen.
(Chapmann-Enskog)

Problème : $\mathbf{c} \in \mathcal{R}^3 \Rightarrow$ Il faut discrétiser l'espace des vitesses pour résoudre numériquement.

Vitesses Discrètes

- 1964 Premier réseaux de vitesses par J. Broadwell puis R. Gatignol.

Gaz sur Réseaux

- 1973 Réseau HPP (Hardy, Pomeau, Pazzis) gaz sur réseau (automates cellulaires)
- 1986 Synthèse des automates cellulaires par S.Wolfram.
- 1988 Premiers réseaux à distributions réelles par McNamara & Zanetti.

Boltzmann sur Réseau

- 1992 D'Humièrre et Qian propose le modèle D2Q9.
- 1992 Modèle à temps de relaxation multiples (d'Humièrre, Lallemand)
- 2002 MRT 3D (d'Humièrre et al.)
- 2006 Cascaded LBM (Geier)
- 2013 Modèle à vitesses relatives (Février, Graille, Dubois)
- 2015 Cumulant Model (Geier)

La discrétisation des vitesses se fait en assurant l'égalité des moments continus et discrets :

- Le nombre de vitesse détermine le nombre de moments conservés.
- L'équilibre est écrit sous forme polynômial.
- Les premiers réseaux compacts conservent les 2 premiers moments. (ρ et ρu). Le moment d'ordre 3 (ρe) possède un terme d'erreur en $\mathcal{O}(Ma^3)$.
- Les réseaux compacts sont donc isothermes et faiblement compressibles.

Lattice Boltzmann Method

Lattice Boltzmann algorithm :

$$g_{\alpha}^{coll}(\mathbf{x}, t) = g_{\alpha}(\mathbf{x}, t) - \frac{1}{\tau_g} [g_{\alpha}(\mathbf{x}, t) - g_{\alpha}^{eq}(\mathbf{x}, t)]$$

MPI communication

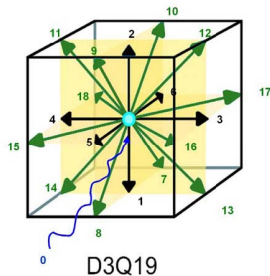
$$g_{\alpha}(\mathbf{x}, t) = g_{\alpha}^{coll}(\mathbf{x} - \mathbf{c}_{\alpha} \Delta t, t - \Delta t)$$

Conditions aux limites

$$\rho = \sum_{\alpha} g_{\alpha}$$

$$\rho \mathbf{u} = \sum_{\alpha} \mathbf{c}_{\alpha} g_{\alpha}$$

$$g_{\alpha}^{eq}(\mathbf{x}, t) = \rho \omega_{\alpha} \left(1 + \frac{\mathbf{u} \cdot \mathbf{c}_{\alpha}}{\tilde{c}_0^2} + \frac{(\mathbf{u} \cdot \mathbf{c}_{\alpha})^2}{2\tilde{c}_0^4} - \frac{|\mathbf{u}|^2}{2\tilde{c}_0^2} \right)$$



$$p = \tilde{c}_0^2 \rho$$

$$\tilde{c}_0^2 = 1/3$$

$$\omega_0 = 1/3$$

$$\omega_{1-6} = 1/18$$

$$\omega_{7-18} = 1/36$$

Multiple relaxation times :

$$\mathbf{m}_\alpha^* = \mathbf{m}_\alpha - S(\mathbf{m}_\alpha - \mathbf{m}_\alpha^{\text{eq}})$$

MPI communication

$$g_\alpha(\mathbf{x} + \mathbf{c}_\alpha \Delta t, t + \Delta t) = M^{-1} \mathbf{m}_\alpha^*(\mathbf{x}, t)$$

Conditions aux limites

$$\rho = \sum_\alpha g_\alpha$$

$$\rho \mathbf{u} = \sum_\alpha \mathbf{c}_\alpha g_\alpha$$

$$g_\alpha^{\text{eq}}(\mathbf{x}, t) = S \mathbf{m}_\alpha^{\text{eq}}$$

$$(m_1, \dots, m_9) = (\rho, e, \epsilon, \rho u_x, q_x, \rho u_y, q_y, p_{xx}, p_{xy})$$

$$S = \text{diag}(0, s_e, s_\epsilon, 0, s_q, 0, s_q, s_\nu, s_\nu) \quad (4)$$

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Stabilité linéaire du modèle de Navier-Stokes isotherme :

$$\frac{\partial \mathbf{U}'}{\partial t} + \mathbf{M}_1 \frac{\partial \mathbf{U}'}{\partial x_1} + \mathbf{M}_2 \frac{\partial \mathbf{U}'}{\partial x_2} + \mathbf{M}_3 \frac{\partial \mathbf{U}'}{\partial x_3} = 0 \quad (5)$$

$$\omega \mathbf{U}' = \mathbf{M}^{\text{NS}} \mathbf{U}' \quad (6)$$

Solution : 3 modes convectifs (Shu & Kovasnay)

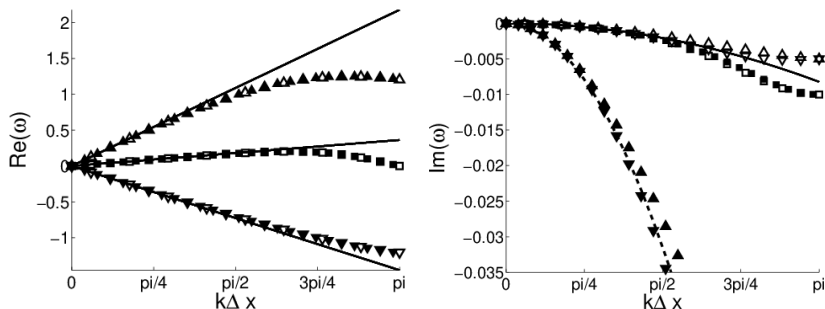
$$\left\{ \begin{array}{lcl} \omega_1 & = & \mathbf{k} \cdot \mathbf{u}_0 - i|\mathbf{k}|^2 \mathcal{N} + |\mathbf{k}|c_0 \sqrt{1 - \left(\frac{|\mathbf{k}|\mathcal{N}}{c_0}\right)^2} \\ \omega_2 & = & \mathbf{k} \cdot \mathbf{u}_0 - i|\mathbf{k}|^2 \mathcal{N} - |\mathbf{k}|c_0 \sqrt{1 - \left(\frac{|\mathbf{k}|\mathcal{N}}{c_0}\right)^2} \\ \omega_3 & = & \mathbf{k} \cdot \mathbf{u}_0 - i|\mathbf{k}|^2 \nu \\ \omega_4 & = & \omega_3 \\ \omega_5 & = & \mathbf{k} \cdot \mathbf{u}_0 \end{array} \right. \quad (7)$$

Stabilité linéaire de la LBM-BGK : Linéarisation de la fonction d'équilibre :

$$f^{eq}(f_{\alpha}^{(0)} + f'_{\alpha}) = f_{\alpha}^{eq,(0)} + \frac{\partial f_{\alpha}^{eq}}{\partial f_{\beta}} \Big|_{f_{\beta}=f_{\beta}^{(0)}} f'_{\alpha} + o(f_{\alpha}'^2) \quad (8)$$

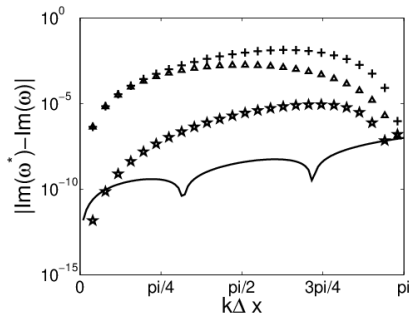
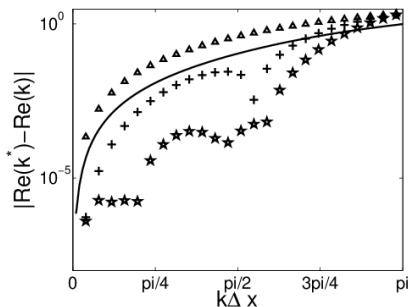
Problème aux valeurs propres :

$$\omega \mathbf{f}' = \mathbf{M}^{\text{LBM}} \mathbf{f}' \quad (9)$$

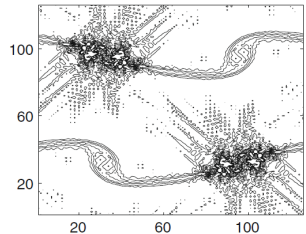
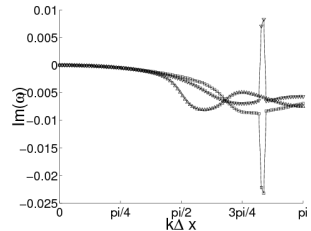


Comparaison avec différences finies :

Espace	Temps	Mach	Mode Acoustique
Ordre 2	3-pas	0.2	$\triangle$
DRP	3-pas	0.2	$+$
Ordre 6 optimisé	6-pas opt.	0.2	$\star$



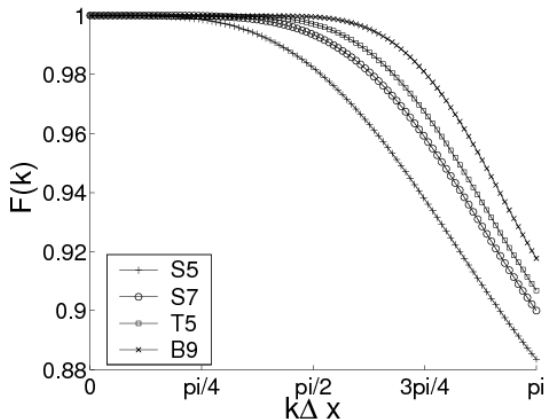
- La dissipation intrinsèque de LBGK est faible.
- Nature compressible du schéma : chaque perturbations se propage de façon isotrope sans être dissipée.
- Problème pour les écoulements haut Reynolds => Instabilités
- Beaucoup de solutions existent pour stabiliser et permettre la simulation des grands voir très grands Reynolds. Beaucoup se font au détriment de la dissipation des ondes acoustiques.
- Application d'un filtrage explicite due à la nature hautes fréquences des instabilités. (Ricot et al. 2009)



Basic spatial filtering of a given quantity Q is performed as follow :

$$\langle Q(\mathbf{x}) \rangle = Q(\mathbf{x}) - \sigma \sum_{j=1}^D \sum_{n=-N}^N d_n Q(\mathbf{x} + n\Delta x_j) \quad (10)$$

$0 < \sigma = Cste < 1$ and coefficients d_n depending on the filter stencil.



Adaptive filtering :

Le rôle du filtre passe-bas est double :

- 1 Supprimer les oscillations hautes fréquences.
- 2 Apporter la dissipation numérique consistante avec un modèle de sous-maille. (Implicit SGS model)

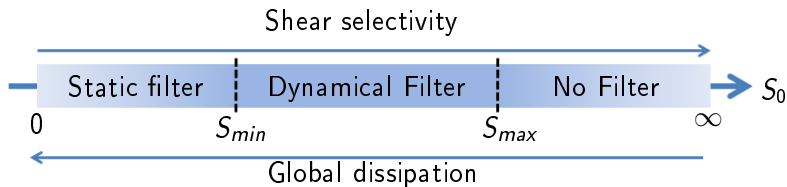
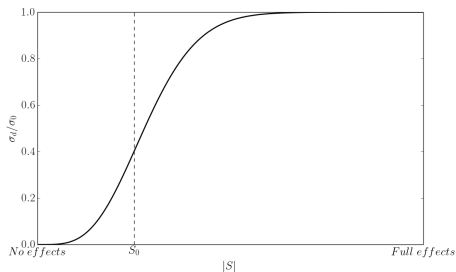
Comme le filtrage est nécessaire surtout dans les zones cisailé, on fait dépendre le coefficient σ du taux de cisaillement local :

$$\sigma_a(\mathbf{x}) = \sigma_0 \xi(|S|) \quad (11)$$

$$\xi(|S|) = \left(1 - e^{-(|S(\mathbf{x})|/S_0)^2}\right)^2 \quad (12)$$

$|S| = \sqrt{2S_{ij}S_{ij}}$, S_0 is a threshold to be determined.

$$\sigma_a(\mathbf{x}) = \sigma_0 \left(1 - e^{-(|S|/S_0)^2} \right)^2$$



$$\sigma_d(\mathbf{x}) = \sigma_0 \left(1 - e^{-(|S|/S_0)^2}\right)^2$$

Threshold estimation :

S_0 could be linked with the maximum amount of shear stress S_{max} :

$$S_0 = \epsilon S_{max}$$

- S_{max} imposed to a constant.
- S_{max} computed from mean field.
- S_{max} based on g_α positivity.

LBM shear computation :

From a LBM framework, the shear stress can be obtained with :

$$2\rho\nu S_{ij} = - \sum_{\alpha} c_{\alpha,i} c_{\alpha,j} (g_{\alpha} - g_{\alpha}^{eq}) \quad (13)$$

Assuming the positivity of g_{α} and $\sum_{\alpha} c_{\alpha,i} c_{\alpha,j} g_{\alpha}^{eq} = \rho u_i u_j$

$$|S| = \frac{Q_f}{2\rho\nu} = \frac{\sqrt{2 \sum_{\alpha} c_{\alpha,i} c_{\alpha,j} g_{\alpha}^{eq} \sum_{\alpha} c_{\alpha,i} c_{\alpha,j} g_{\alpha}^{eq}}}{2\rho\nu} \leq \frac{\sqrt{2} u_i u_j}{2\nu}$$

Then :

$$S_{max} = \frac{\sqrt{2} U_0^2}{2\nu} \sim \frac{Re_{\delta} U_0}{\delta}$$

Choosing the filtered quantity

3 different possibilities with increasing computational cost :

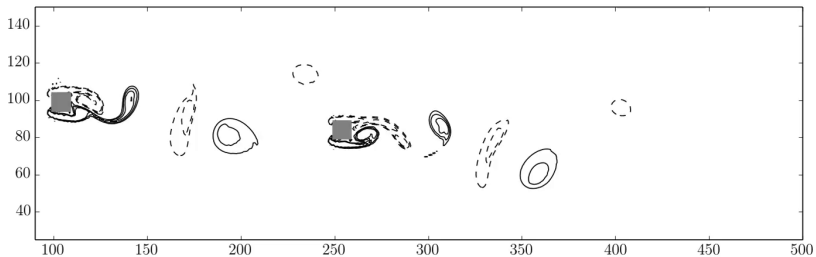
- 1 Filtering momenta : ρ, u_x, u_y, u_z (4 tables).
- 2 Filtering distribution functions : g_α (19 tables).
- 3 Filtering collision operator : $-\frac{1}{\tau_g}(g_\alpha - g_\alpha^{eq})$ (19 tables).

Then the overall algorithm becomes :

Modified algorithm for moments filtering :

- Estimation of S_0 .
- Compute $|S|$ with (13) and update σ_d
- Collision Step
- (Collision Filtering)
- Propagation Step
- Update moments
- (Moments filtering)
- Update g_α^{eq}

An example of adaptivity :



Taylor Green Vortex

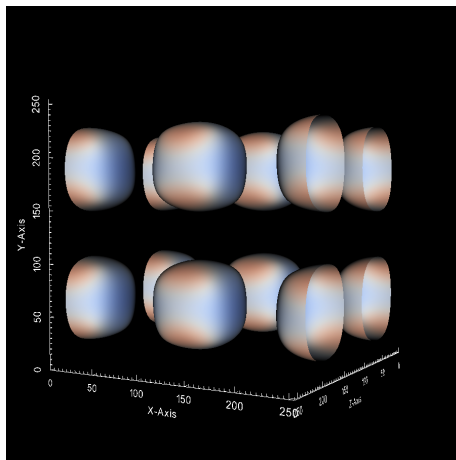
Initialization of the Taylor-Green vortex is done by setting velocity and pressure in a cubic and fully periodic domain of size 2π :

$$\begin{cases} p &= p_{\infty} + \frac{\rho_{\infty} U_{\infty}^2}{16} [\cos(2z) + 2][\cos(2x) + \cos(2y)] \\ u &= U_{\infty} \sin(x) \cos(y) \cos(z) \\ v &= -U_{\infty} \cos(x) \sin(y) \cos(z) \\ w &= 0 \end{cases} \quad (14)$$

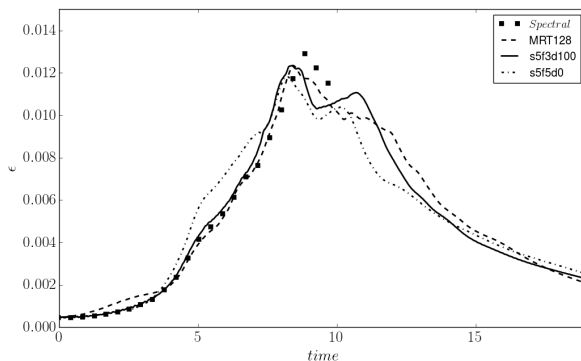
Parameters :

- Imposing $Re = 1600$, $M_{\infty} = 0.085$, $\rho_{\infty} = 1$, $dx = 2\pi/n_x$
- Induces : $U_{\infty} = 0.049$, $p_{\infty} = 1/3$, $\tau_g = 0.5 + 1.46n_x 10^{-5}$

Taylor-Green vortex at $Re = 1600$:



Dissipation rate of kinetic energy

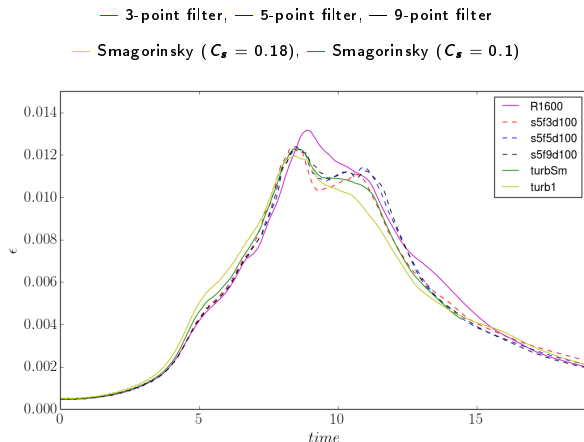


Comparison between Adaptive and classical techniques.

- Results on a 128^3 grid
- MRT results from P.Lallemand
- Adaptive filtering give similar results than the MRT model.
- Adaptive filtering with 3-point stencil gives better results than static 5-point.

Comparison to subgrid scale model

- Adaptive filtering gives better results than SGS model.
- SGS describes poorly the laminar-turbulent transition region.
- SGS strategy gives similar results as adaptive filtering with a low shear stress sensitivity S_0 .



Comparison between adaptive filtering and subgrid-scale model

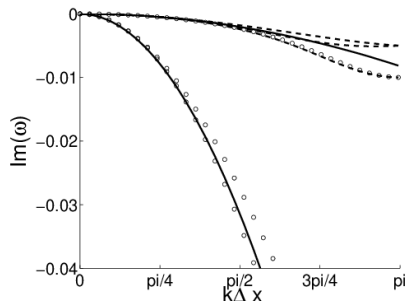
Adaptive Relaxation Times (ART)

MRT Collision

$$\mathbf{m}_\alpha^* = \mathbf{m}_\alpha - S(\mathbf{m}_\alpha - \mathbf{m}_\alpha^{eq})$$

$$g_\alpha(\mathbf{x} + \mathbf{c}_\alpha \Delta t, t + \Delta t) = M^{-1} \mathbf{m}_\alpha^*(\mathbf{x}, t)$$

In the MRT framework, the filtering step is not needed.



$$(m_1, \dots, m_9) = (\rho, e, \epsilon, \rho u_x, q_x, \rho u_y, q_y, p_{xx}, p_{xy})$$

The function $\xi(|S|)$ is used to switch from MRT in sheared region to BGK in uniform flow region :

$$S' = \text{diag} (0, s'_e, s'_\epsilon, 0, s'_q, 0, s'_q, s'_\nu, s'_\nu) \quad (15)$$

where $s' = s + (1 - \xi)(s_\nu - s)$

Then the overall algorithm becomes :

Modified algorithm for adaptive relaxation times :

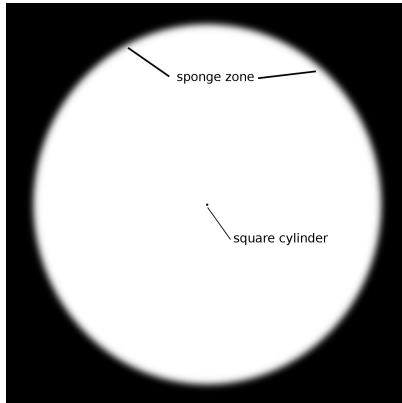
- Compute $|S|$ with (13) and update ξ
- Update relaxation time matrix S'
- Collision Step in the momentum space
- Propagation Step
- Update moments
- Update equilibrium

Sound radiated by a square cylinder

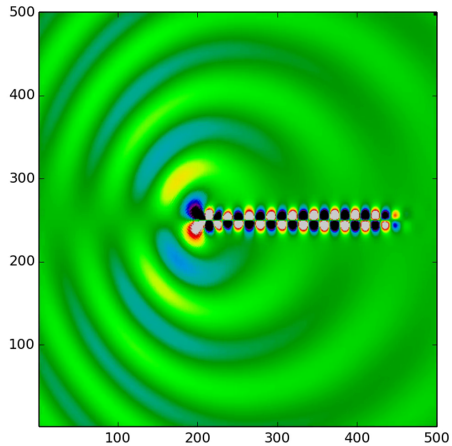
Parameters

- Reynolds $R_D = 500$
- Mach number $M_a = 0.3$
- the square size is $D = 5$
- the domain size is $300D \times 300D$
- The domain is initialized with a uniform flow and periodic boundary conditions.
- A circular sponge zone is defined at the domain boundary to damp the outgoing structures.

Case setup



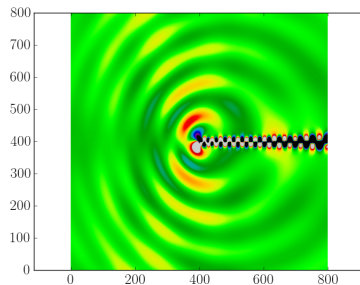
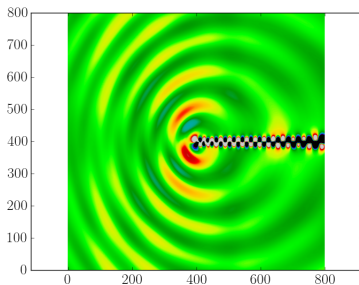
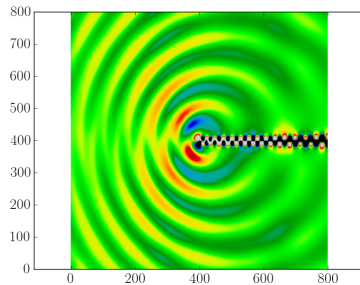
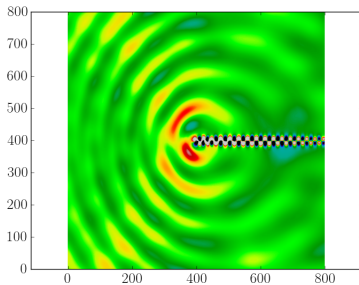
- Details about LBM sponge zone from Hui et al 2011.
- Here the filter coefficient σ and the relaxation time τ_ν are set to high values in the sponge zone.
- For ART model, the switcher is set to 1 in the sponge zone.



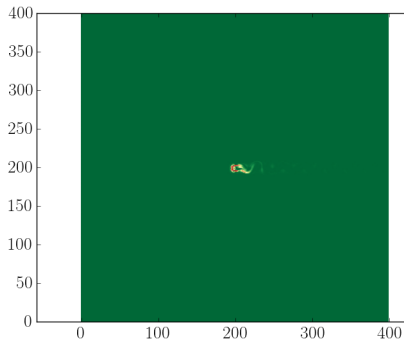
BGK

MRT

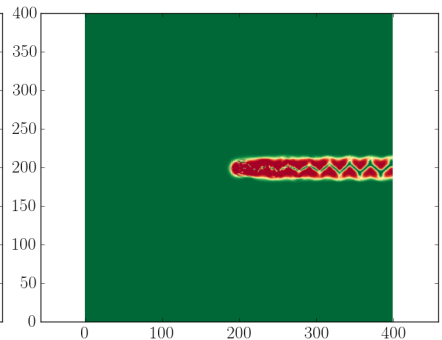
Standard

Adaptive
 $S_0 = 1.0$ 

Maps of $\xi(|S|)$



$S_0 = 0.1$



$S_0 = 1.0$

Perspectives

- Traitement de paroi optimisé pour les modèles adaptatifs.
- Etude d'un couplage NS/LBM
- Etude des modèles Cascaded et Cumulant
- Etude d'un cas acoustique hautes fréquences.
- Comparaison avec les approches lagrangiennes.

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He-Chen-Zhang model

1 additional distribution function to track index function ϕ :

BGK on a D2Q9 lattice with force terms :

$$f_{\alpha}^{+} = \text{BGK}(f_{\alpha}) - \frac{2\tau - 1}{2\tau} (\mathbf{c}_{\alpha} - \mathbf{u}) \cdot \frac{\nabla \psi(\phi)}{RT} \Gamma_{\alpha}(\mathbf{u})$$

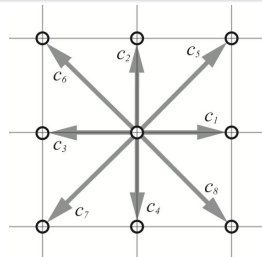
$$g_{\alpha}^{+} = \text{BGK}(g_{\alpha}) - \frac{2\tau - 1}{2\tau} (\mathbf{c}_{\alpha} - \mathbf{u}) \cdot [\nabla \psi(\rho)(\Gamma_{\alpha}(\mathbf{u}) - \Gamma_{\alpha}(0)) - \Gamma_{\alpha}(\mathbf{u})(\mathbf{F}_s + \mathbf{G})]$$

Equilibrium state :

$$\Gamma_{\alpha}(\mathbf{u}) = \omega_{\alpha} \left(1 + \frac{\mathbf{u} \cdot \mathbf{c}_{\alpha}}{RT} + \frac{(\mathbf{u} \cdot \mathbf{c}_{\alpha})^2}{2(RT)^2} - \frac{|\mathbf{u}|^2}{2RT} \right)$$

$$f_{\alpha}^{eq} = \phi \Gamma_{\alpha}(\mathbf{u})$$

$$g_{\alpha}^{eq} = \omega_{\alpha} \left[p + \rho RT \left(\frac{\mathbf{u} \cdot \mathbf{c}_{\alpha}}{RT} + \frac{(\mathbf{u} \cdot \mathbf{c}_{\alpha})^2}{2(RT)^2} - \frac{|\mathbf{u}|^2}{2RT} \right) \right]$$



He-Chen-Zhang model

Moments :

$$\phi = \sum_{\alpha} f_{\alpha}$$

$$p_h = \sum_{\alpha} g_{\alpha} - \frac{1}{2} \mathbf{u} \cdot \nabla \psi(\rho)$$

$$\rho RT \mathbf{u} = \sum_{\alpha} \mathbf{c}_{\alpha} g_{\alpha} + \frac{RT}{2} (\mathbf{F}_s + \mathbf{G})$$

Thermodynamics :

$$\psi(\rho) = p_h - \rho RT$$

$$\psi(\phi) = p_{th} - \phi RT$$

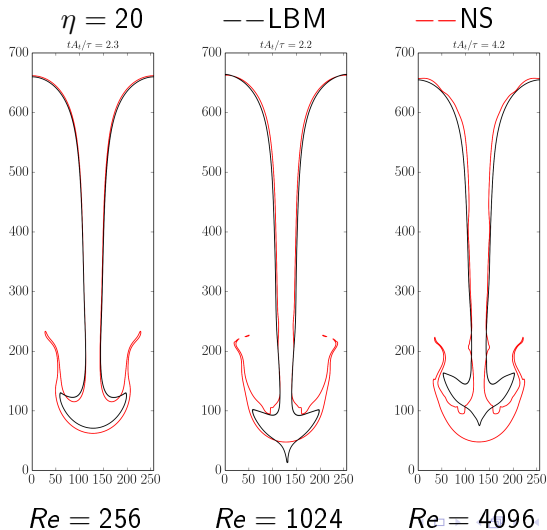
Equation d'état : Carnahan-Starling

$$p_{th} = \phi RT \frac{1 + \phi + \phi^2 - \phi^3}{(1 - \phi)^3} - a\phi^2$$

Défauts de stabilité dans les zones à fort gradient de densité $\Rightarrow$ Utilisation du filtrage adaptatif basé sur $\nabla \rho$.

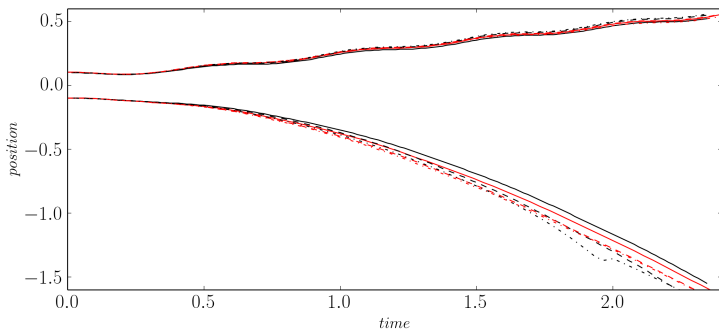
Rayleigh-Taylor Instability $A_t = 0.9$

NS Results from V.Daru (2nd order TVD + SuperBee limiter)

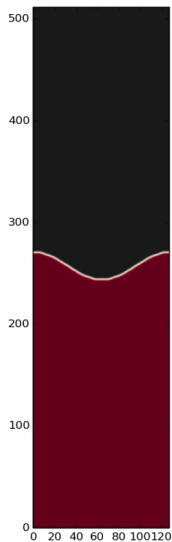
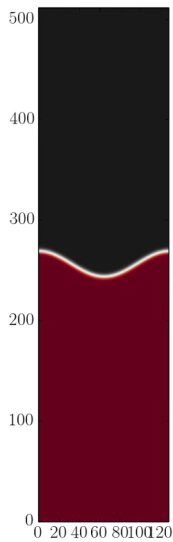
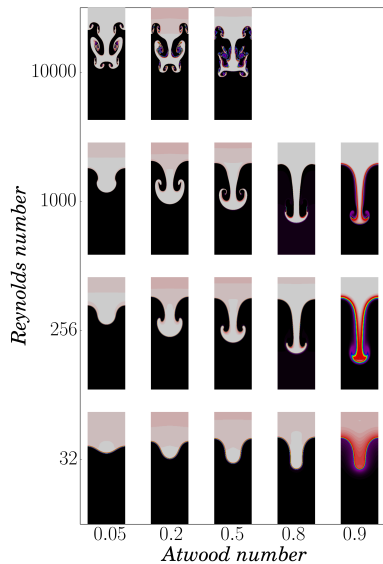


$$A_t = 0.9$$

Spike and bubble position.



— $Re = 256$, — — $Re = 1024$, ... $Re = 4096$


 $A_t = 0.5$

 $A_t = 0.9$


- Work from PhD student L.Vienne ending in 2019.
- Build a LBM model for mixing simulation of miscible gas.
- Independent of the collision operator.
- Taking diffusion into account through forcing term.

Multicomponent LBM :

$$f_{\alpha}^m(\mathbf{x} + \mathbf{e}_{\alpha}\delta_t, t + \delta_t) = f_{\alpha}^m(\mathbf{x}, t) - \frac{1}{\tau_m} \left[f_{\alpha}^m(\mathbf{x}, t) - f_{\alpha}^{m(eq)}(\mathbf{x}, t) \right] + S_{\alpha}^m(\mathbf{x}, t)$$

Forcing term :

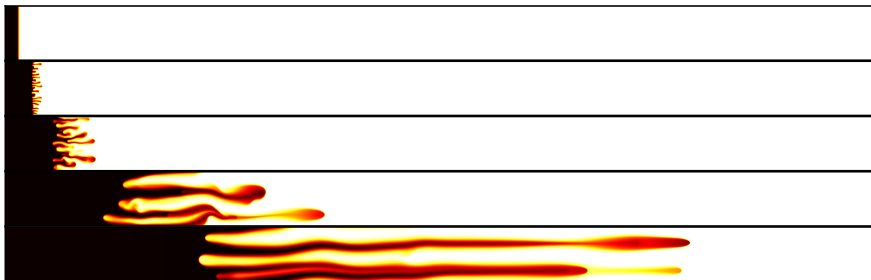
$$S_{\alpha}^m = (1 - \frac{\delta_t}{2\tau_m})\omega_{\alpha} \left[\frac{\mathbf{e}_{\alpha} \cdot \mathbf{u}_m}{c_s^2} + \frac{(\mathbf{e}_{\alpha} \cdot \mathbf{u}_m)\mathbf{e}_{\alpha}}{c_s^4} \right] \cdot \mathcal{F}_m$$

Diffusion term between species : Kerkhof & Geboers

$$\mathcal{F}_m = -p \sum_{n=1}^N \frac{x_m x_n}{\mathcal{D}_{mn}} (\mathbf{u}_m - \mathbf{u}_n),$$

Application to the viscous fingering instability (Saffman-Taylor)

Porous medium is modeled with Gray Lattice Boltzmann (I.Ginzburg).



Perspectives

- Etude de l'instabilité de ST à l'échelle du pore.
- Etude d'une couche de mélange transitoire multiespèces.
- Caractérisation numérique de la diffusion dans les approches LBM-multiespèces.
- Etudes des faibles nombres de Schmidt.